

Introduction of a 3D Particle-In-Cell (PIC) method in a Discontinuous Galerkin electromagnetic model.

Application to high power microwave sources (NADEGE project)

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Vlasov-Maxwell equations

Maxwell equations

- Find $(\mathbf{E}, \mathbf{H}) : \Omega \times [0, T] \rightarrow \mathbb{R}^3 \times \mathbb{R}^3$ such that :

$$\begin{cases} \varepsilon \frac{\partial \mathbf{E}}{\partial t} - \nabla \times \mathbf{H} + \mathbf{J} = 0 \text{ on } \Omega \\ \mu \frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0 \text{ on } \Omega \\ \mathbf{E} \times \mathbf{n}(\mathbf{x}) = 0 \text{ on } \partial \Omega \\ \nabla \cdot \mathbf{D} = \rho \\ \nabla \cdot \mathbf{B} = 0 \\ \mathbf{E}(\mathbf{x}, 0) = \mathbf{E}_0(\mathbf{x}) \text{ and } \mathbf{H}(\mathbf{x}, 0) = \mathbf{H}_0(\mathbf{x}) \text{ on } \Omega \end{cases} \quad (1)$$

- Three-dimensional computational domain Ω .
- \mathbf{E} is the electric field and \mathbf{H} the magnetic field with \mathbf{E}_0 and \mathbf{H}_0 initial conditions.
- ε , μ and σ denote, respectively, the permittivity, the permeability and the conductivity of the medium.
- $\mathbf{J}(\mathbf{x}, t)$ and $\rho(\mathbf{x}, t)$ represent electric current and charge densities.

Vlasov equation

- Evolution of the particle distribution function f in a collisionless plasma :

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \nabla f + \frac{q}{m} (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial \mathbf{f}}{\partial \mathbf{v}} = 0 \quad (2)$$

This equation is written for each species of particles.

- Invariance of this distribution in time along trajectories submitted to electromagnetic fields.

- Positions and velocities of particles (\mathbf{x}, \mathbf{v}) are solutions of the characteristics of the Vlasov equation (motion equations) :

$$\begin{cases} \frac{d\mathbf{x}}{dt} = \mathbf{v} \\ \frac{d\mathbf{v}}{dt} = \frac{q}{m} [\mathbf{E} + \mathbf{v} \wedge \mathbf{B}] \end{cases} \quad (3)$$

Coupling of equations

- Charge conservation equation :

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{J} = 0 \quad (4)$$

- Numerical problem : the discrete charge conservation equation is not satisfied.

- Solution : introduction of a corrector term on electric field \mathbf{E} at each time step.

- Choice of a purely hyperbolic correction (appropriate for our schemes)

$$\begin{cases} \varepsilon \frac{\partial \mathbf{E}}{\partial t} - \nabla \times \mathbf{H} + \mathbf{J} + \chi \nabla \rho = 0 \\ \mu \frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0 \\ \mu \frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{E} = \chi \frac{\rho}{\varepsilon_0} \\ \nabla \cdot \mathbf{B} = 0 \end{cases} \quad (5)$$

Discontinuous Galerkin formulation

- Three-dimensional computational domain $\Omega = \cup_{i=1}^N K_i$, $(K_i)_{i=1,\dots,N}$ hexahedra.
- On each cell K_i , Maxwell equations with hyperbolic correction are rewritten :

$$\begin{cases} \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t} - \nabla \times \mathbf{H} + \chi \nabla \rho + \mathbf{J} + \alpha [\mathbf{H} \times \mathbf{n}] + \beta [\mathbf{n} \times (\mathbf{E} \times \mathbf{n})] + \tau \chi [\mathbf{n} \cdot \mathbf{on}] = 0 \\ \mu_0 \frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} + \delta [\mathbf{E} \times \mathbf{n}] + \gamma [\mathbf{n} \times (\mathbf{H} \times \mathbf{n})] = 0 \\ \mu_0 \frac{\partial \rho}{\partial t} - \chi \frac{\rho}{\varepsilon_0} + \chi \nabla \cdot \mathbf{E} + \eta \chi [\mathbf{E} \cdot \mathbf{n}] - \theta [\mathbf{o}] = 0 \end{cases} \quad (6)$$

- Conditions : $1 + \alpha - \gamma = 0$, $\beta, \delta \geq 0$ and $1 + \tau + \eta = 0$.

- Boundary conditions for the corrective terms :

$$\begin{aligned} \eta &= -1, \theta = 1, \tau = 0 \text{ (absorbing boundary conditions)} \\ \eta &= \theta = 0, \tau = -1 \text{ (metallic boundary conditions)} \end{aligned} \quad (7)$$

- $(\mathbf{E}, \mathbf{H}) \in (\mathcal{V}_r)^2$, $\mathcal{V}_r = \{\mathbf{v} \in [L^2(\Omega)]^3 : \forall K \in T, DF_K^* \mathbf{v}|_K \circ F_K \in [Q_r(K)]^3\}$.
- Elements of the mesh : a conform mapping F_K such that $F_K(K) = K$.
- For a spatial approximation order r , $(r+1)^3$ degrees of freedom located at the Gauss quadrature points and 3 polynomial basis functions φ at each degree.
- Degrees of freedom and basis functions φ , on each element K : $\varphi \circ F_K(\hat{\mathbf{x}}) = DF_K^{-1} \varphi(\hat{\mathbf{x}})$, with DF_K the Jacobian matrix.

Particle-In-Cell method

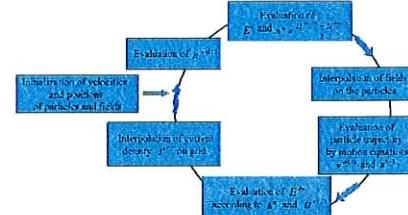


FIGURE 1: Principle of 3D PIC method.

- Current density \mathbf{J} and charge density ρ spawn by charged particles :

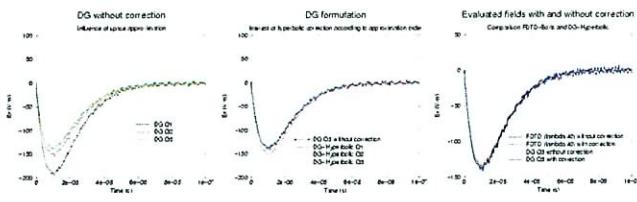
$$\begin{aligned} \rho(\mathbf{x}, t) &= \sum_s q_s \int_{\mathbb{R}^3} f_s(\mathbf{x}, \mathbf{v}, t) d\mathbf{p} \\ \mathbf{J}_s(\mathbf{x}, t) &= \sum_s q_s \int_{\mathbb{R}^3} \mathbf{p} f_s(\mathbf{x}, \mathbf{v}, t) d\mathbf{p} \end{aligned} \quad (8)$$

- Interpolation of electromagnetic fields on the particles and particles on mesh with Nearest Grid Point (NGP) method.

Numerical results

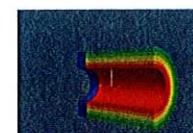
Comparison : DG-Hyperbolic and FDTD-Boris

- Boris' method (elliptic-hyperbolic correction)
- Geometry : perfectly metallic cubic cavity
- Introduction of current and charge densities in the cavity



Coaxial line: $R_{ext} = 0.2m$, $R_{int} = 0.1m$, $L = 0.6m$ and $V_{peak} = 4MV$

- Source : Transverse ElectroMagnetic (TEM) mode



- Field emission (breakdown field = $2.5 \times 10^7 \text{ V/m}$)

